How is it then that many protons are stuck to each other (and to neutrons) inside an atomic nucleus? Because the "strong force" acting only along very short distances (~ 10<sup>-15</sup> m) attracts them to each other (by a factor of 100 compared to electric repul sion) There is also a "weak force" responsible for radioctivity § 8.1 Coulomb's law How do we measure the electric charge of? What is the force between two static charges q, and q, as a function of their positions?

-s answers are contained in Coulomb's low" (Charles-Augustine de Coulomb (1736 - 1806))

$$\overline{F_{12}} = -\overline{F_{21}} = \frac{q_1 q_2}{4\pi \epsilon_0 r_{12}} \overline{e}_{12}$$
where  $r_{12} = |\overline{r_2} - \overline{r_1}|$  and  $\overline{e}_{12} = \frac{\overline{r_0} - \overline{r_1}}{|\overline{r_2} - \overline{r_1}|}$ 

$$\overline{F_2} = \frac{q_1 \overline{r_2} - \overline{r_1}}{r_1} = \frac{q_1}{q_2} \frac{\overline{r_2} - \overline{r_1}}{r_2}$$
Unit of charge : 1 Coulomb [C]  
 $\frac{1}{4\pi \epsilon_0} = q_1 \cdot 10^9 \frac{N \cdot m^2}{C^2}$ 
If two charges, 1 coulomb each, are  
separated by 1 meter, the repulsive force  
between them will be  $q \cdot 10^7 N$ .  
(weight of about 10,000,000 adults)  
In these units the charge of the proton,  
denoted by c, is 1.6 \cdot 10^{-19} C and that  
of the electron is -e.

Properties of charge: · charge is conserved "locally"; suppose charge would be conserved non-locally:  $\begin{vmatrix} & & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\$ -> due to special relativity, only simultaneous in one reference frame ( can choose a reference frame where there are at an instance two particles or no particle) -> conservation laws that hold in all frames must be local charge is quantized : All charges we have ever seen are integral multiples of a basic unit of charge, e= 1.6.10<sup>-19</sup>C

Superposition principle Consider 3 charges q1, q2, and q3 at F, F and F3 F<sub>13</sub> 9 7 1 F<sub>23</sub> 'J\_ F31 €  $\vec{F}_{31} + \vec{F}_{32} = \vec{F}_{3(1+1)} + \vec{F}_{3(1+1)}$ What will be the force on from the other two? answer:  $\overline{F}_{3(1+2)} = \overline{F}_{31} + \overline{F}_{31}$ " superposition principle" (encountered already in Newtonian mechanics) Question: Suppose we want to verify Coulomb's law. How will you confirm the dependence of the force on q1, q2, and r=1r,-r,1? Thik about it.

The ratio of gravitational to electric forces  
Claim: Fg/Fe, ratio of gravitational  
to electric forces, ~ 10-40  
T  
How do we obtain  
this answer?  
Consider two particles of masses m,  
and m, and charges q, and q, a  
distance r apart.  

$$\rightarrow \frac{Fq}{Fe} = \left(\frac{Gm_{1}m_{2}}{r^{2}}\right) \cdot \left(\frac{q_{1}q_{2}}{4\pi \varepsilon_{r}c^{2}}\right)^{-1} = \frac{Gm_{1}m_{2}\cdot 4\pi \varepsilon}{q_{1}q_{2}}$$
  
For two electrons, this gives:  
 $\frac{Fq}{Fe} = \frac{6\cdot7\cdot10^{-11}\cdot(q\cdot10^{-31})^{2}}{(1.6\cdot10^{19})^{2}\cdot q\cdot10^{5}} = 2\cdot3\cdot10^{-43}$   
proton and electron:  $\frac{Fq}{Fe} \sim 10^{-36}$   
But: in maeroscopic bodies, electric charges  
cane in pairs and cancel!  
That's why the earth and the  
moon attract each other by gravity!

Coulomb's law for continuous charge density Example 1: Consider a circular wire of radius R with & Coulombs per meter, lying in xy-plane dF 7 R Find the electric force it exerts on a point charge q'located at height z on the z-axis.

Solution:  

$$I dFI = \frac{q \cdot \lambda dl}{4\pi \cdot \varepsilon_0 (R^2 + \cdot \cdot \cdot^2)}$$
Only need keep the vertical part  
of the force, the hovizontal part  
is cancelled by  $dl^*$ :  
 $F_2 = \int dF_2 = \int \frac{q \cdot \lambda dl}{4\pi \cdot \varepsilon_0 (R^2 + \cdot \cdot^2)} \times \frac{2}{\sqrt{R^2 + \cdot^2}}$   
 $= \frac{q \cdot \lambda - 2\pi \cdot R \cdot 2}{4\pi \cdot \varepsilon_0 (R^2 + \cdot \cdot^2)^{3/2}}$ 
where the factor  $\frac{2}{\sqrt{R^2 + \cdot^2}} = \cos \theta$   
projects out the vertical part of dF  
and  $2\pi \cdot R$  is the integral over dl.  
Tests:  
i) set  $z=0 \longrightarrow$  should give zero  
by symmetry  
and indeed if does!  
2) for  $2 \gg \cdot R$  the wire should  
appear like a point charge, it does!

$$\frac{\S 8.2 \quad \text{The electric field } \vec{E}}{\texttt{Xet} \quad us \quad write \quad \text{the force } q_1 \, due \ \text{to } q_1 : \\ \vec{F}_{21} = \frac{q_1 q_2}{4\pi \epsilon_0 r_1^2} \vec{e}_{12}$$

$$= \frac{q_1}{4\pi \epsilon_0 r_1^2} \vec{e}_{12} \cdot q_2 = \vec{E}(\vec{r}_2) q_2$$

$$\vec{E}(\vec{r}_2) : \quad \text{'the electric field at the location of } q_2''$$

$$\frac{\text{Step 1: The charge } q_1 \text{ produces a field}}{\vec{E}(\vec{r}_2) = \frac{q_1}{4\pi \epsilon_0 r_1^2} \vec{e}_{12} \text{ at the location of } q_2}$$

$$\frac{\text{Step 2: The charge } q_1 \text{ responds to the } field \quad \text{by feeling a force } \vec{F}_1 = q_1 \vec{E}(\vec{v}_2) = \frac{F_1}{4\pi \epsilon_0 r_1^2} \vec{e}_{12} \text{ at the location of } q_2$$

$$\frac{\text{Step 2: The charge } q_1 \text{ responds to the } field \quad \text{by feeling a force } \vec{F}_1 = q_1 \vec{E}(\vec{v}_2) = \frac{F_1}{4\pi \epsilon_0 r_1^2} \vec{e}_{12} \vec{e}_{12} \vec{e}_{13} \vec{e}_{1$$

Visualizing the field  
For a point-charge at the origin,  
the electric field is:  

$$\vec{E}(\vec{r}) = \frac{q}{4\pi \xi r^2} \vec{e_r}$$
  
 $= \frac{q}{4\pi \xi r^3} \vec{r}$  field lines  
 $\vec{r}$  field lines

Question. Do we lose information dout the electric field at q poit F if we connect arrows to lines? 1) We retain the information about the "direction" of the electric field, i.e. the direction of the force vector acting on a "probe charge" 2) It seems that we are losing information about the "magnitude of the force But: "density" of field lines through a fixed area DA contains information about magnitude! area of a sphere in IR3 grows as r2 with the radius r - density of N field lines at distance r:  $\sim \frac{N}{r^2} \sim |\vec{E}|$ 

Proposition 1: The electric field of a point-charge is a conservative field. That is, the force on a probe charge is a conservative force. Proof. According to Theorem 3 in \$4, a force is conservative if it can be written as  $\vec{F} = \vec{\nabla} U(\vec{r}) = \begin{pmatrix} \partial_x U \\ \partial_y U \\ \partial_z U \end{pmatrix}$   $\rightarrow for the electric field$   $\vec{E}$  we have to show:  $\vec{E} = -\vec{\nabla} V$  $\frac{\text{(laim: V(r)} = \frac{q}{4\pi \varepsilon_r}}{4\pi \varepsilon_r}$ We compute.  $-\frac{\partial V}{\partial x} = -\frac{q}{4\pi \varepsilon} \frac{\partial (1/r)}{\partial x} = -\frac{q}{4\pi \varepsilon} \frac{\partial}{\partial x} \left[ \frac{1}{\sqrt{x^2 + y^2 + z^2}} \right]$  $= \frac{q}{4\pi\xi_{0}} \left[ \frac{1}{2} \right] \frac{2x}{(x^{2}+y^{2}+z^{3})^{3}/2} = \frac{q}{4\pi\xi_{0}} \frac{1}{z^{2}} \frac{x}{z}$  $= \overline{\nabla} V = \frac{q}{4\pi\epsilon_{r}} \frac{\overline{e_{x}} \times + \overline{e_{y}} \cdot q}{r} = \overline{E}(\overline{r})$ 



 $\vec{E}_{+} = \frac{q}{4\pi\epsilon_{0}} \frac{\vec{r}_{+}}{r_{+}^{3}} = \frac{q}{4\pi\epsilon_{0}} \frac{\vec{e}_{1}(x-a) + \vec{e}_{2} \cdot q}{((x-a)^{2} + q^{2})^{3} \cdot 2}$   $\vec{F}_{-} = \frac{q}{4} \vec{r}_{-} \qquad q = \vec{e}_{1}(x+q) \cdot \vec{e}_{1} \cdot q$ 

 $\vec{E}_{-} = -\frac{q}{4\pi\epsilon_{o}} \frac{\vec{r}_{-}}{r^{3}} = -\frac{q}{4\pi\epsilon_{o}} \frac{\vec{e}_{i}(x+a) + \vec{e}_{2}q}{(x+a)^{2} + q^{2})^{2}/2}$   $= \sum_{-} \vec{E}_{-} = \frac{q}{4\pi\epsilon_{o}} \left[ \frac{\vec{e}_{i}(x-a) + \vec{e}_{2}q}{(x-a) + \vec{e}_{2}q} - \frac{\vec{e}_{i}(x+a) + \vec{e}_{2}q}{(x+a)^{2} + q^{2})^{2}/2} \right]$ 

For a point A with  $x \ge a$ , we obtain:  $\overline{E} = \frac{q}{4\pi \varepsilon_0} \overline{e_1} \left( \frac{1}{(x-a)^2} - \frac{1}{(x+a)^2} \right)$ 

